

# Advanced Topics in Condensed Matter

## Lecture 5: Inelastic scattering

Dr. Ivan Zaluzhnyy

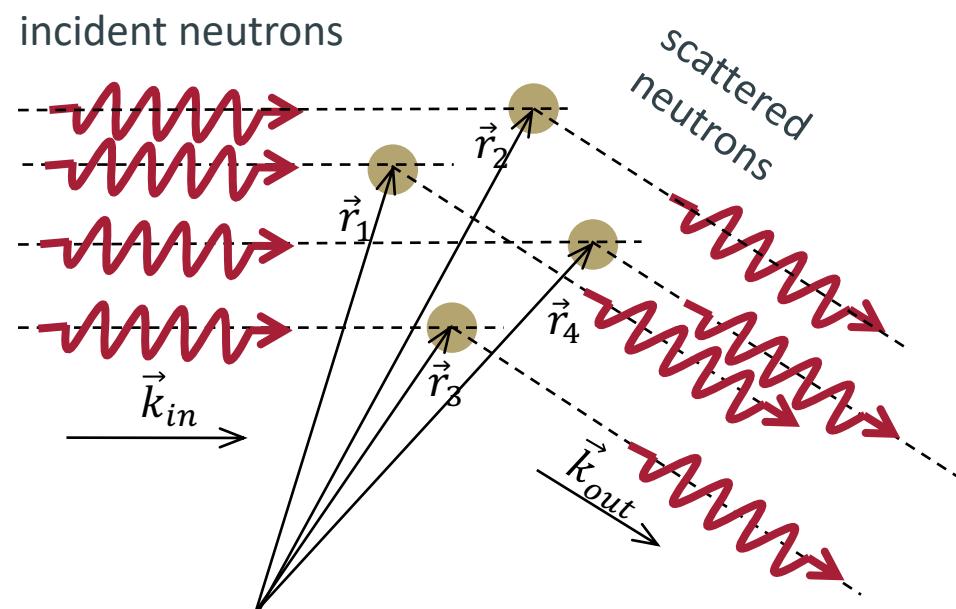
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# Reminder

$$I(\vec{q}) \propto \frac{\sigma_{coh}}{4\pi} N \cdot S(\vec{q}) + \frac{\sigma_{inc}}{4\pi} N \quad - \text{scattered intensity}$$



$$S(\vec{q}) = \frac{1}{N} \left\langle \sum_{n,m=1}^N e^{-i\vec{q}(\vec{r}_n - \vec{r}_m)} \right\rangle$$

- static structure factor

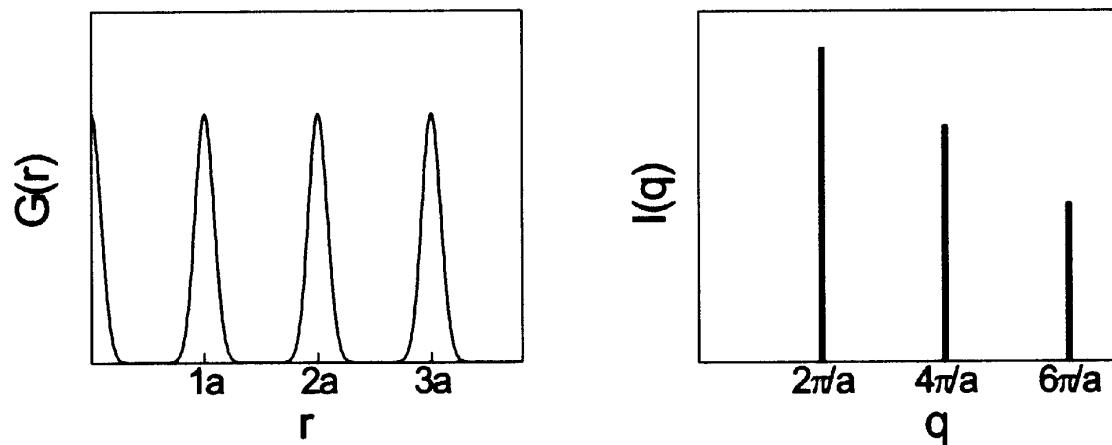
$$S(\vec{q}) = \int G(\vec{r}) e^{-i\vec{q}\vec{r}} d\vec{r}$$

$$G(\vec{r}) = \frac{1}{N} \left\langle \sum_{n,m=1}^N \delta(r - ((\vec{r}_n - \vec{r}_m))) \right\rangle$$

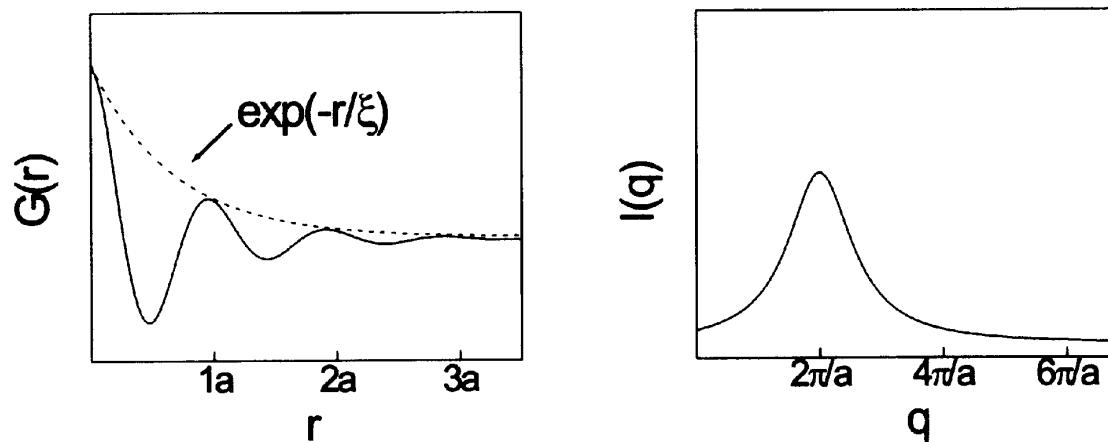
- correlation function

# Correlation function $G(r)$ and structure factor $S(q)$

long-range order (LRO)

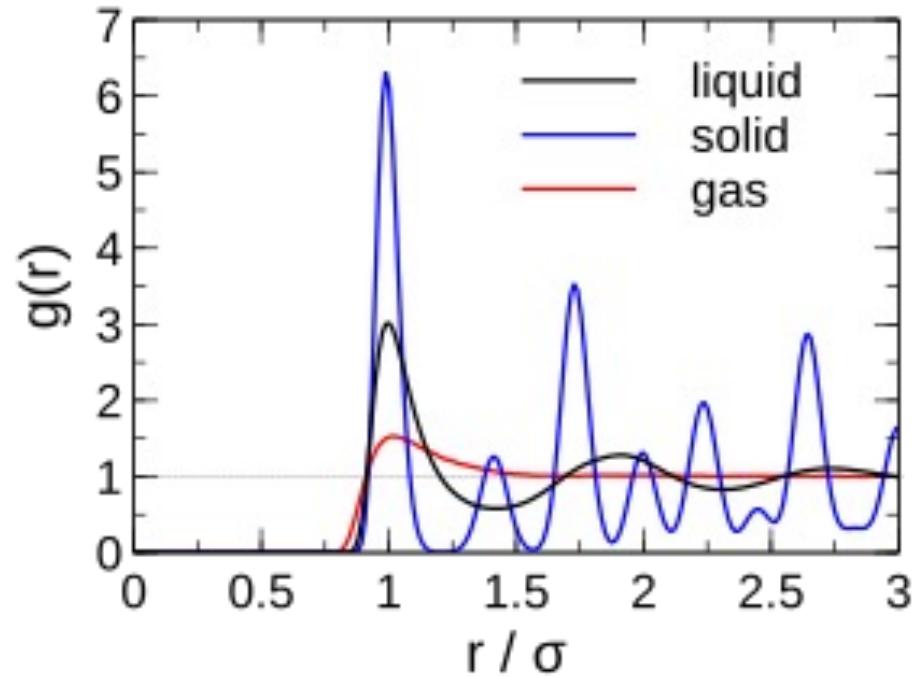


short-range order (SRO)



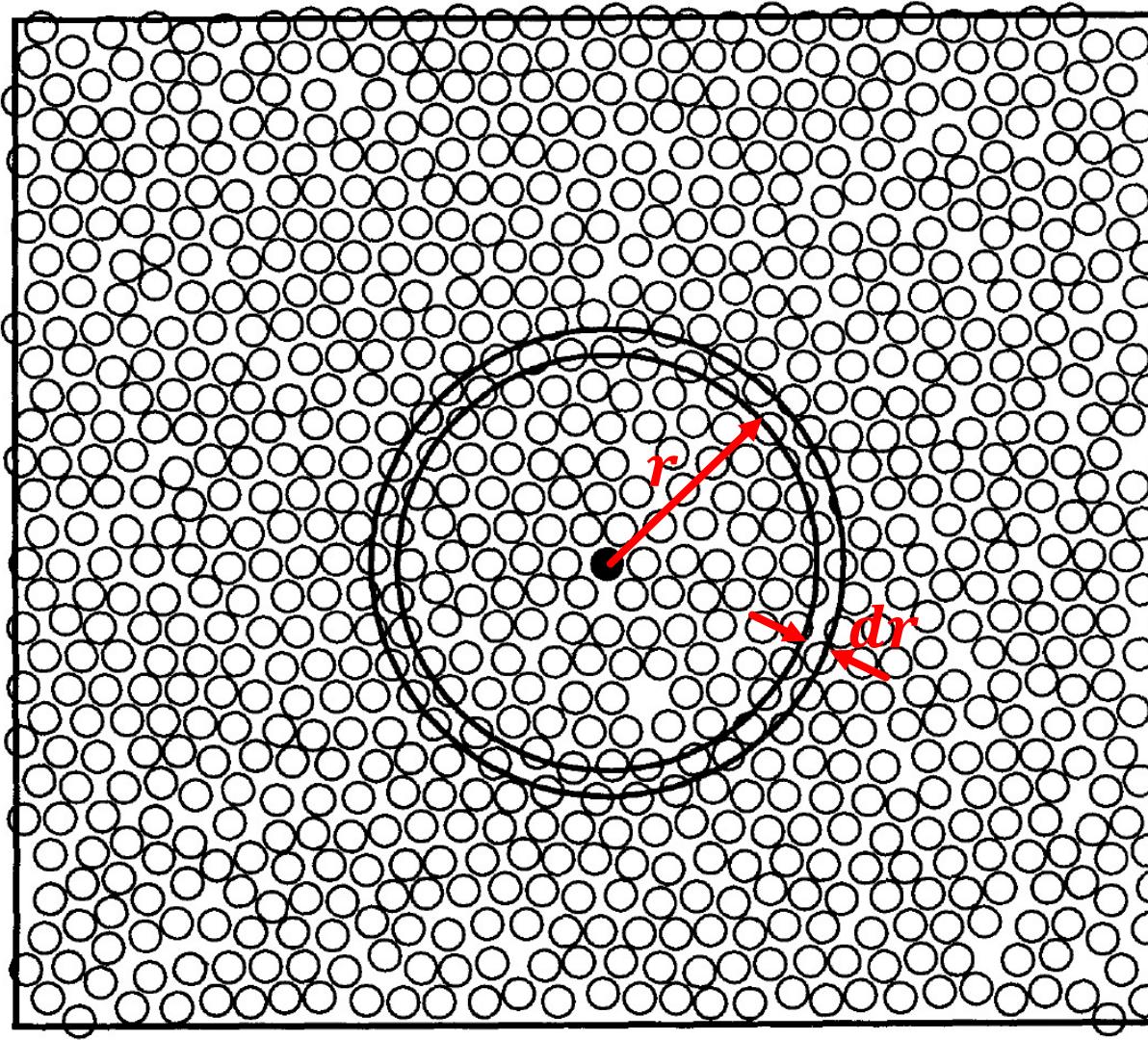
# Pair Distribution Function

$$\begin{aligned}
 G(\vec{r}) &= \frac{1}{N} \left\langle \sum_{\alpha\alpha'} \delta(\vec{r} - (\vec{r}_{\alpha'} - \vec{r}_\alpha)) \right\rangle = \frac{1}{N} \left\langle \sum_{\alpha=\alpha'} \delta(\vec{r} - (\vec{r}_{\alpha'} - \vec{r}_\alpha)) + \sum_{\alpha \neq \alpha'} \delta(\vec{r} - (\vec{r}_{\alpha'} - \vec{r}_\alpha)) \right\rangle = \\
 &= \frac{1}{N} \left\langle N\delta(\vec{r}) + \sum_{\alpha \neq \alpha'} \delta(\vec{r} - (\vec{r}_{\alpha'} - \vec{r}_\alpha)) \right\rangle = \\
 &= \delta(\vec{r}) + \left\langle \frac{1}{N} \sum_{\alpha \neq \alpha'} \delta(\vec{r} - (\vec{r}_{\alpha'} - \vec{r}_\alpha)) \right\rangle = \\
 &= \delta(\vec{r}) + \left\langle \frac{N}{N} \sum_{\alpha' \neq 0} \delta(\vec{r} - (\vec{r}_{\alpha'} - \vec{r}_0)) \right\rangle = \\
 &= \delta(\vec{r}) + \left\langle \sum_{\alpha \neq 0} \delta(\vec{r} - (\vec{r}_\alpha - \vec{r}_0)) \right\rangle = \\
 &= \delta(\vec{r}) + \langle n \rangle g(\vec{r}), \quad \text{where} \\
 \text{PDF} \quad g(\vec{r}) &= \frac{1}{\langle n \rangle} \left\langle \sum_{\alpha \neq 0} \delta(\vec{r} - (\vec{r}_\alpha - \vec{r}_0)) \right\rangle
 \end{aligned}$$



$\langle n \rangle g(\vec{r}) d\vec{r}$  is an average number of particles in the volume  $d\vec{r}$  at the separation  $\vec{r}$  from a given one (excluding self-correlation).

# Pair distribution function

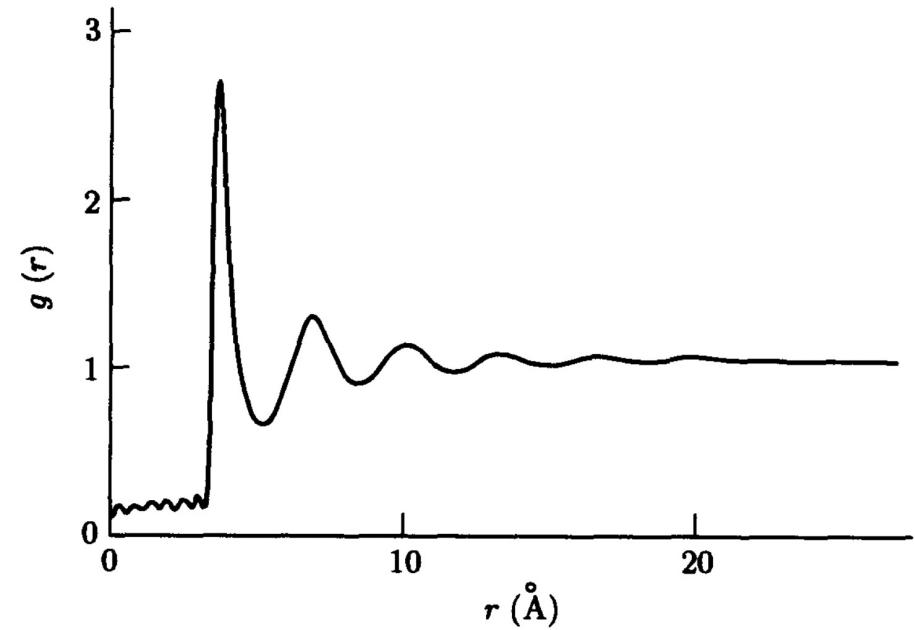
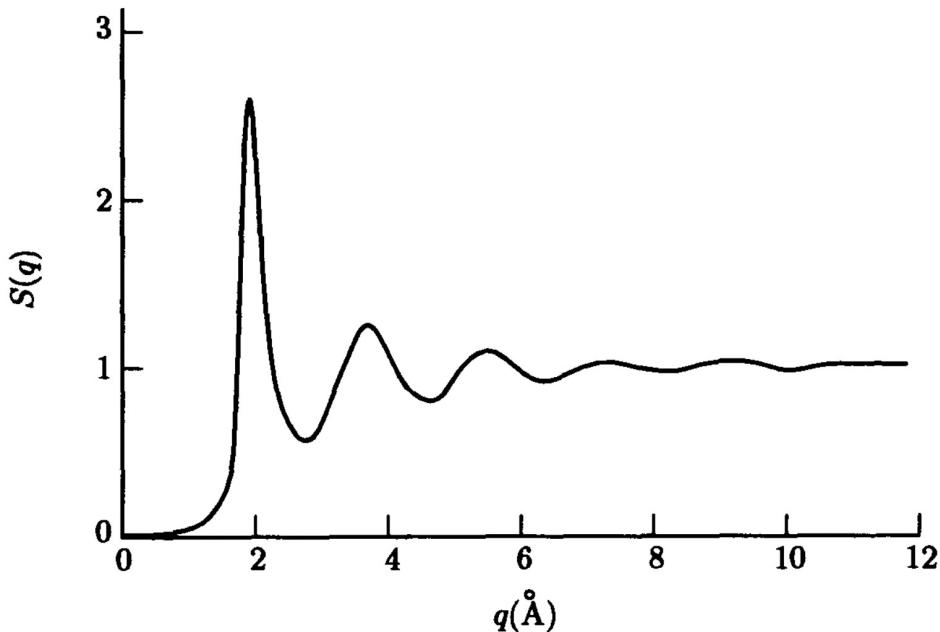


$\langle n \rangle g(r)dr$  – average amount of particles which can be found at the distances from  $r$  to  $r + dr$  from any given particle

P. Chaikin & T. Lubensky (1995)

# Pair distribution function

The structure factor and radial distribution function for liquid argon at T=85 K.



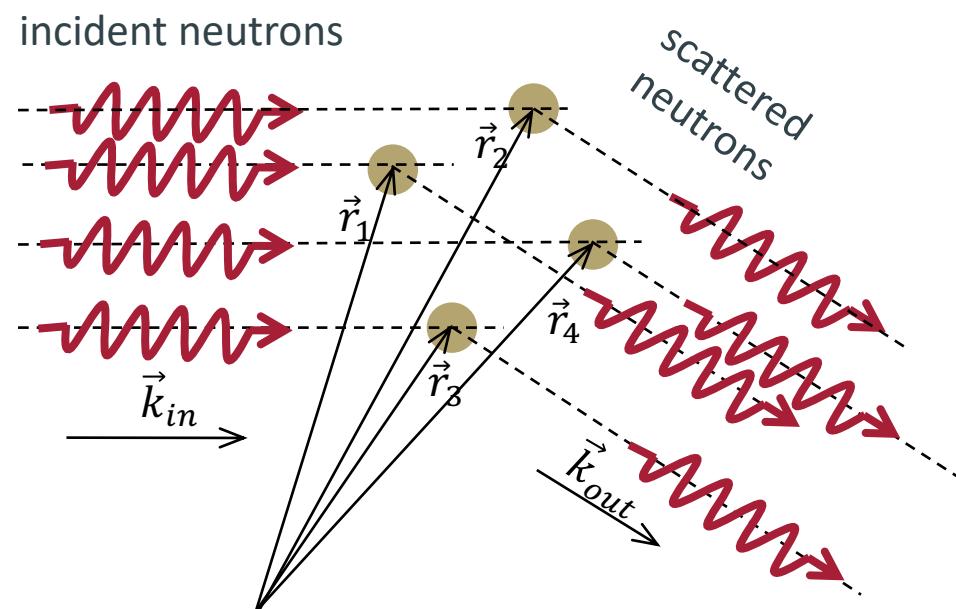
$$g(r) = 1 + \frac{1}{2\pi^2 \langle n \rangle r} \int_0^{+\infty} (S(q) - 1) q \sin(qr) dq$$

P. Chaikin & T. Lubensky (1995)

# Static system of scatterers

$$I(\vec{q}) \propto \frac{\sigma_{coh}}{4\pi} N \cdot S(\vec{q}) + \frac{\sigma_{inc}}{4\pi} N$$

- scattered intensity



$$\vec{q} = \vec{k}_{out} - \vec{k}_{in}$$

$$S(\vec{q}) = \frac{1}{N} \left\langle \sum_{n,m=1}^N e^{-i\vec{q}(\vec{r}_n - \vec{r}_m)} \right\rangle$$

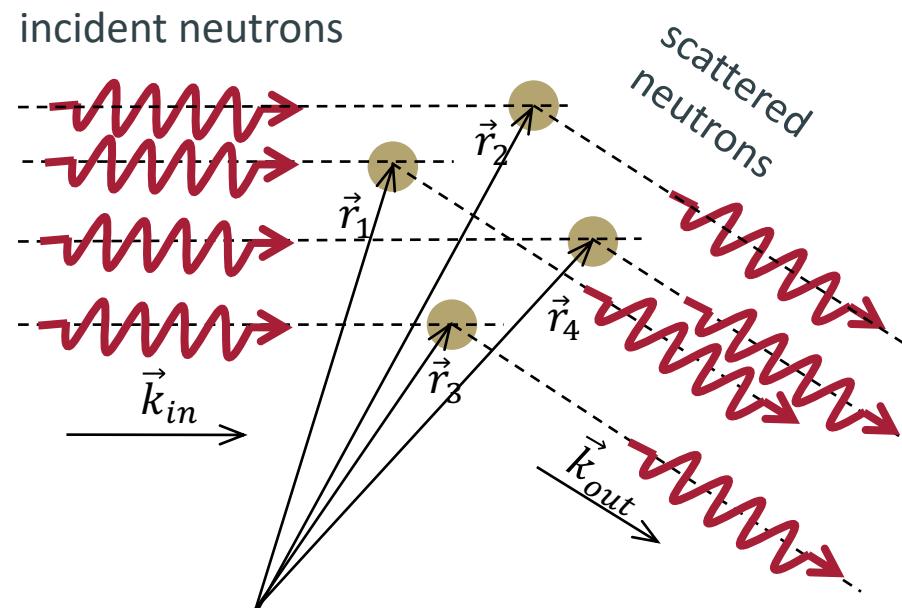
- static structure factor

$$S(\vec{q}) = \int G(\vec{r}) e^{-i\vec{q}\vec{r}} d\vec{r}$$

$$G(\vec{r}) = \frac{1}{N} \left\langle \sum_{n,m=1}^N \delta(r - ((\vec{r}_n - \vec{r}_m))) \right\rangle$$

- correlation function

# System of moving scatterers



$$\vec{q} = \vec{k}_{out} - \vec{k}_{in}$$

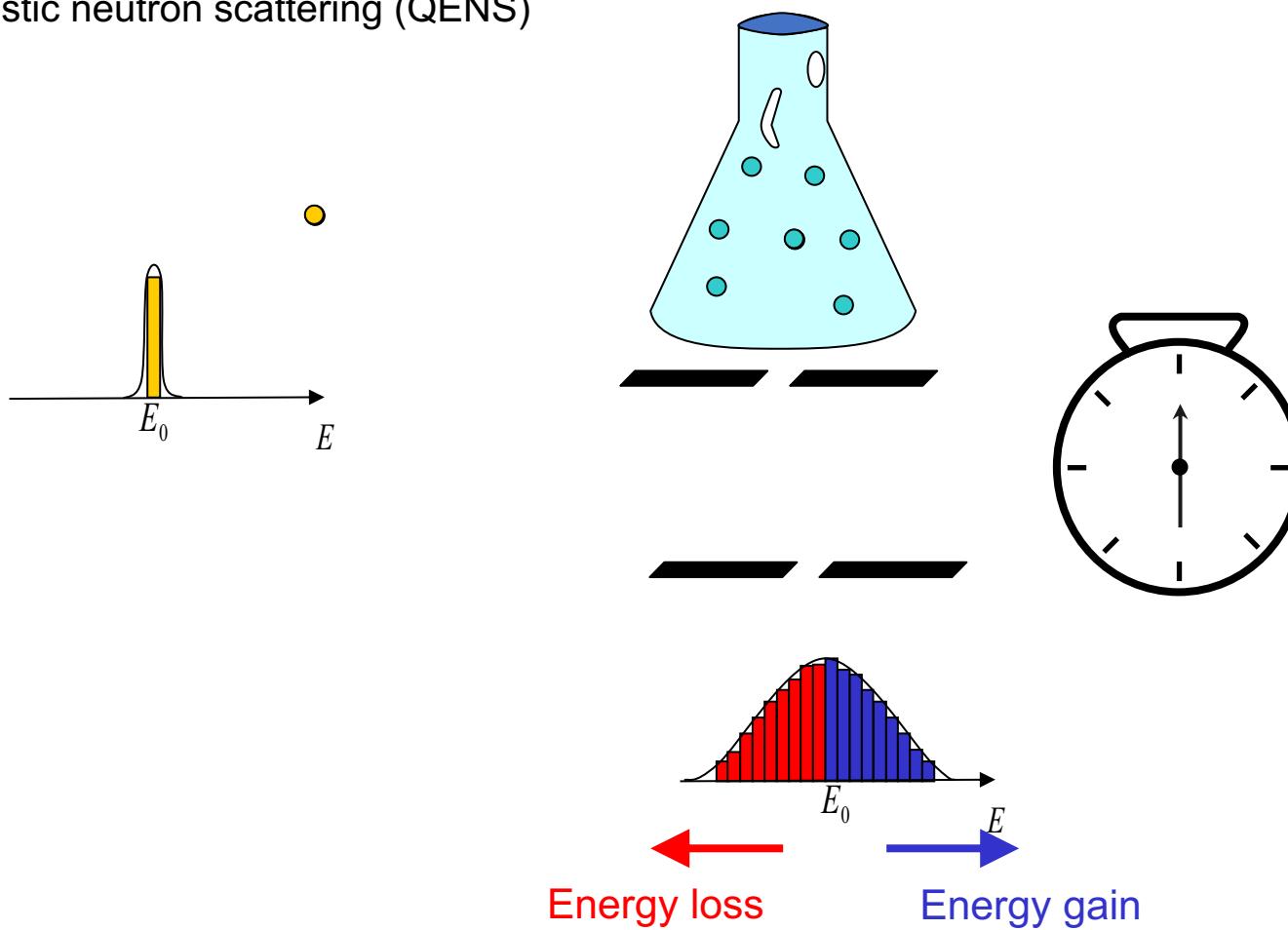
Scatterers are moving, so we have to correlate the positions of the particle "n" at  $t = 0$  with the particle "m" at  $t \neq 0$

$$G(\vec{r}, t) = \frac{1}{N} \left\langle \sum_{n,m=1}^N \delta \left( r - ((\vec{r}_n(t) - \vec{r}_m(0))) \right) \right\rangle$$

- generalized correlation function

# System of moving scatterers

Quasielastic neutron scattering (QENS)



The scattered neutrons will have the energy  $E = E_0 + \hbar\omega$

# Static vs. dynamic

Scatterers have fixed positions:

Scatterers can move:

$$I(\vec{q}) \propto \frac{\sigma_{coh}}{4\pi} N \cdot S(\vec{q}) + \frac{\sigma_{inc}}{4\pi} N$$

$$S(\vec{q}) = \frac{1}{N} \left\langle \sum_{n,m=1}^N e^{-i\vec{q}(\vec{r}_n - \vec{r}_m)} \right\rangle$$

$$S(\vec{q}) = \int G(\vec{r}) e^{-i\vec{q}\vec{r}} d\vec{r}$$

$$G(\vec{r}) = \frac{1}{N} \left\langle \sum_{n,m=1}^N \delta(r - ((\vec{r}_n - \vec{r}_m))) \right\rangle$$

# Static vs. dynamic

Scatterers have fixed positions:

$$I(\vec{q}) \propto \frac{\sigma_{coh}}{4\pi} N \cdot S(\vec{q})$$

$$S(\vec{q}) = \frac{1}{N} \left\langle \sum_{n,m=1}^N e^{-i\vec{q}(\vec{r}_n - \vec{r}_m)} \right\rangle$$

$$S(\vec{q}) = \int G(\vec{r}) e^{-i\vec{q}\vec{r}} d\vec{r}$$

$$G(\vec{r}) = \frac{1}{N} \left\langle \sum_{n,m=1}^N \delta\left(r - ((\vec{r}_n - \vec{r}_m))\right) \right\rangle$$

Scatterers can move:

$$I(\vec{q}, \omega) \propto \frac{k_{out}}{k_{in}} \cdot \frac{\sigma_{coh}}{4\pi} N \cdot S_{coh}(\vec{q}, \omega)$$

$$J(\vec{q}, t) = \frac{1}{N} \left\langle \sum_{n,m=1}^N e^{-i\vec{q}(\vec{r}_n(t) - \vec{r}_m(0))} \right\rangle$$

$$S_{coh}(\vec{q}, \omega) = \frac{1}{2\pi\hbar} \int G(\vec{r}, t) e^{-i\vec{q}\vec{r}} e^{i\omega t} d\vec{r} dt$$

$$G(\vec{r}, t) = \frac{1}{N} \left\langle \sum_{n,m=1}^N \delta\left(r - ((\vec{r}_n(t) - \vec{r}_m(0)))\right) \right\rangle$$

$$S_{coh}(\vec{q}, \omega) = \frac{1}{2\pi\hbar} \int G(\vec{r}, t) e^{-i\vec{q}\cdot\vec{r}} e^{i\omega t} d\vec{r} dt$$

Structure factor

Correlation function

$$S(\vec{q}, \omega)$$

$$\int S(\vec{q}, \omega) d\omega$$

$$\int S(\vec{q}, \omega) d\vec{q}$$

$$\int S(\vec{q}, \omega) d\omega d\vec{q}$$

$$S(\vec{q}, \omega = 0)$$

$$S(\vec{q} = 0, \omega)$$

$$G(\vec{r}, t)$$

$$G(\vec{r}, t = 0)$$

$$G(\vec{r} = 0, t)$$

$$G(\vec{r} = 0, t = 0)$$

$$\int G(\vec{r}, t) dt$$

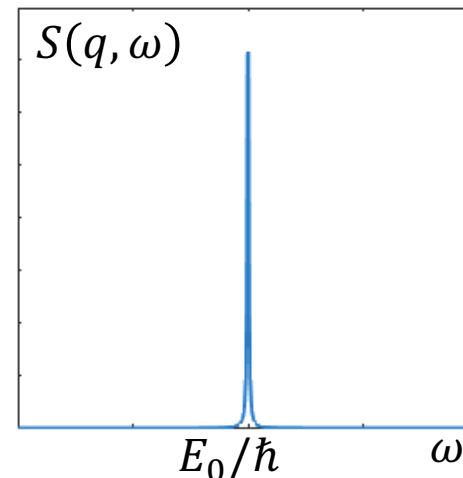
$$\int G(\vec{r}, t) d\vec{r}$$

# Quasielastic neutron scattering

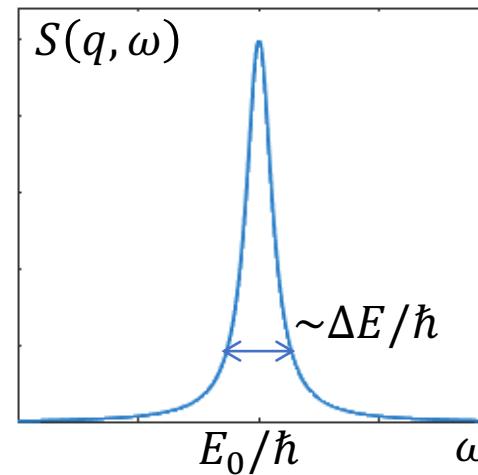
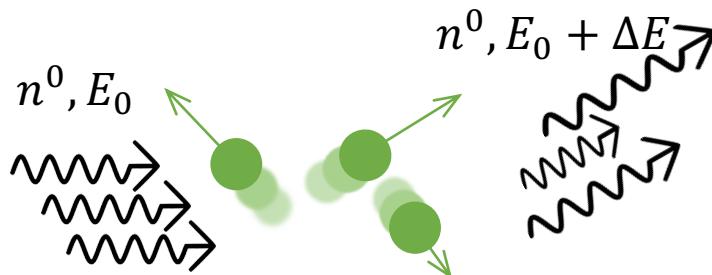
Static sample



Quasielastic neutron scattering  
(spatially incoherent, inelastic)



Dynamic sample



$$\tau \propto 1/\Delta E$$

$$1\text{ps} \Leftrightarrow 4\text{meV}$$

# Van Hove function

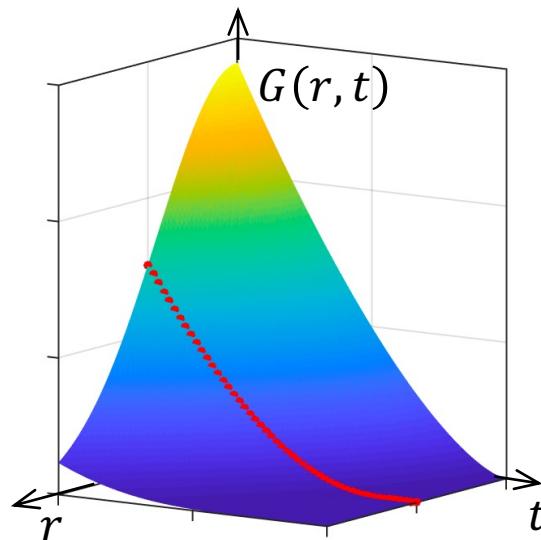
Correlation function in real time and space:  $G(\mathbf{r}, t) = (2\pi)^3 \hbar \int d^3 \mathbf{q} d\omega e^{-i(\mathbf{q}\mathbf{r} - \omega t)} S(\mathbf{q}, \omega)$

Scattering function in frequency domain  
and reciprocal space:  $S(\mathbf{q}, \omega) = \frac{1}{2\pi\hbar} \int d^3 \mathbf{r} dt e^{i(\mathbf{q}\mathbf{r} - \omega t)} G(\mathbf{r}, t)$

$$D\nabla^2 n(\mathbf{r}, t) = \frac{\partial}{\partial t} n(\mathbf{r}, t)$$

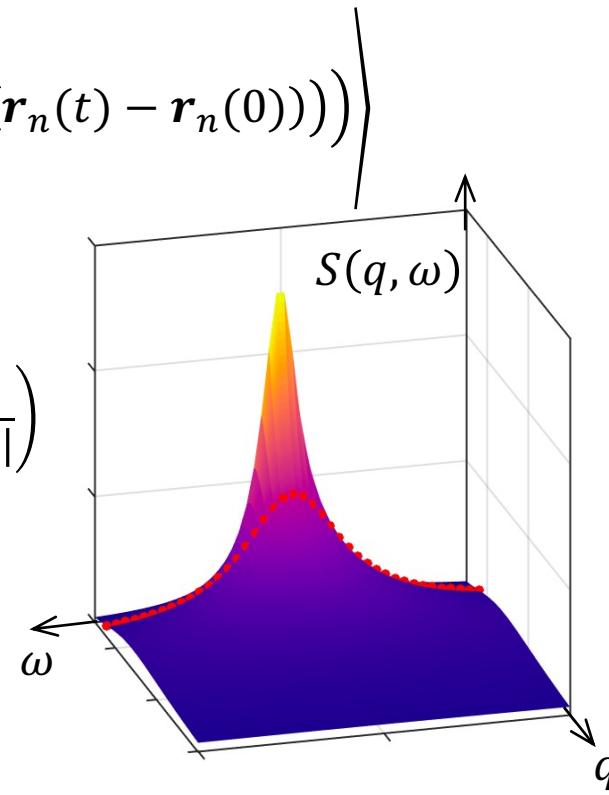
$$G(\mathbf{r}, t) = \frac{1}{N} \langle n(\mathbf{r}, t) n(0, 0) \rangle = \frac{1}{N} \left\langle \sum_{n=1}^N \delta \left( \mathbf{r} - ((\mathbf{r}_n(t) - \mathbf{r}_n(0))) \right) \right\rangle$$

Diffusion:

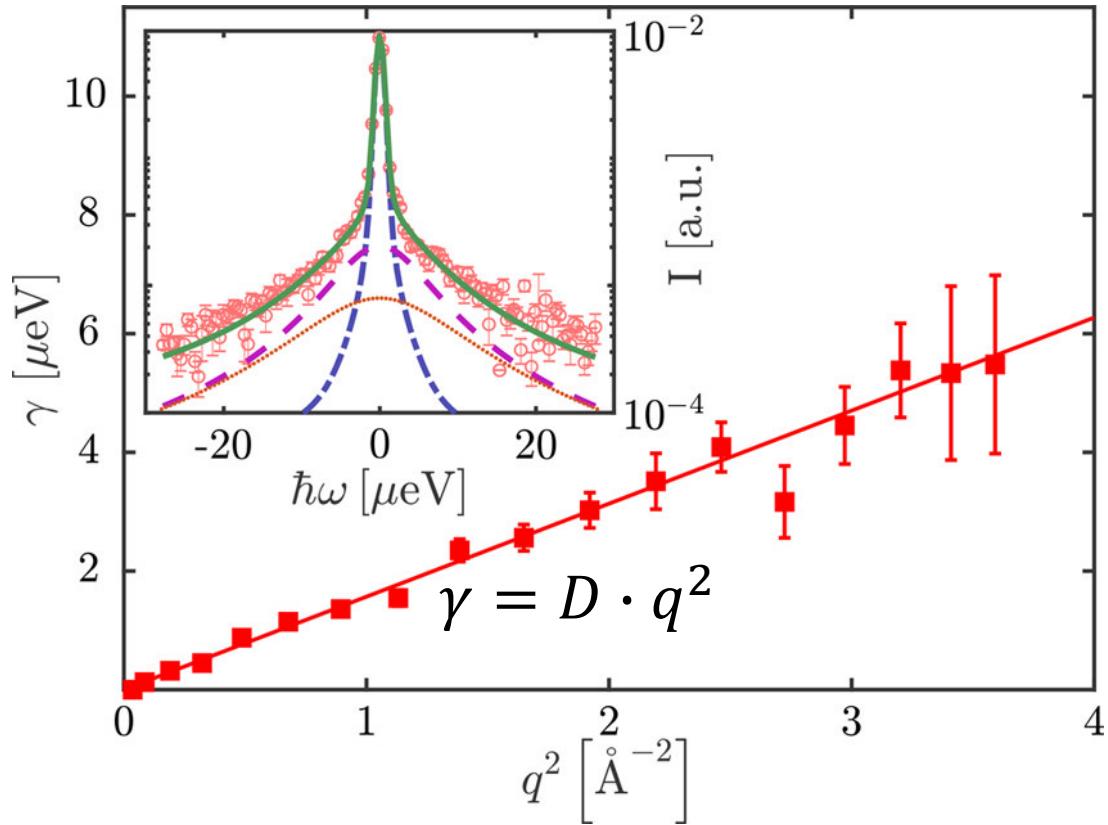
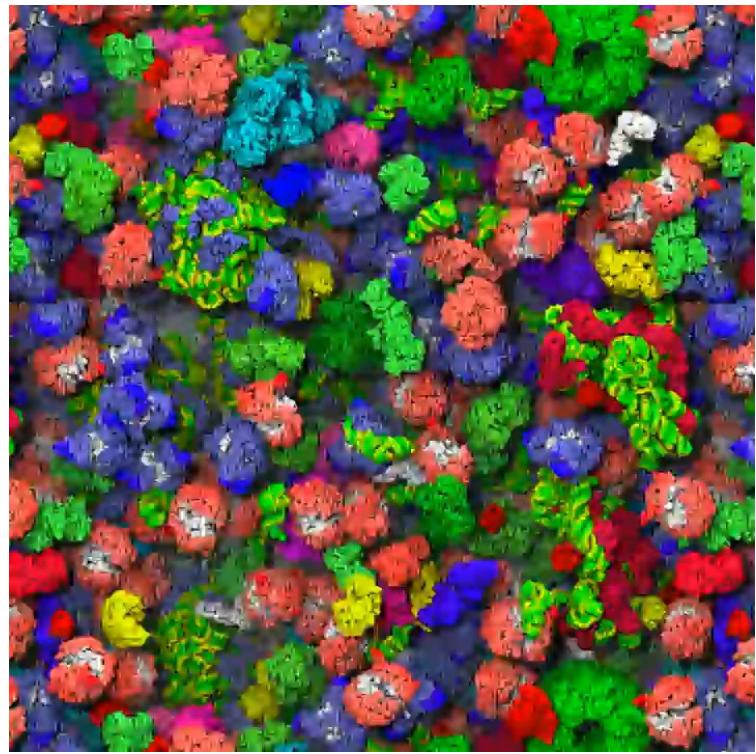


$$G(r, t) = \frac{1}{(4\pi D|t|)^{3/2}} \exp \left( -\frac{r^2}{4D|t|} \right)$$

$$S(q, \omega) = \frac{1}{\pi} \frac{Dq^2}{(Dq^2)^2 + \omega^2}$$



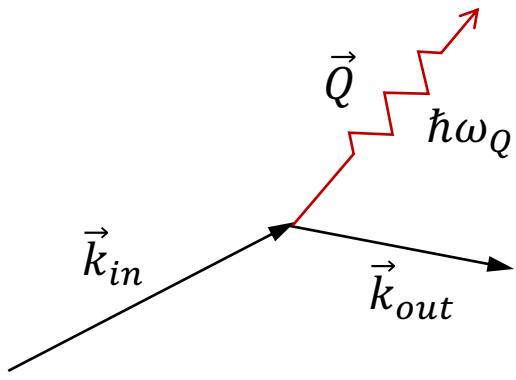
# Measuring self-diffusion



McGuffee & Elcock, PLOS Comp. Biol. 6(3): e1000694 (2010)

M. Grimaldo et al., J. Phys. Chem. Lett. 10 (2019), 1709

# Scattering of neutrons by phonons

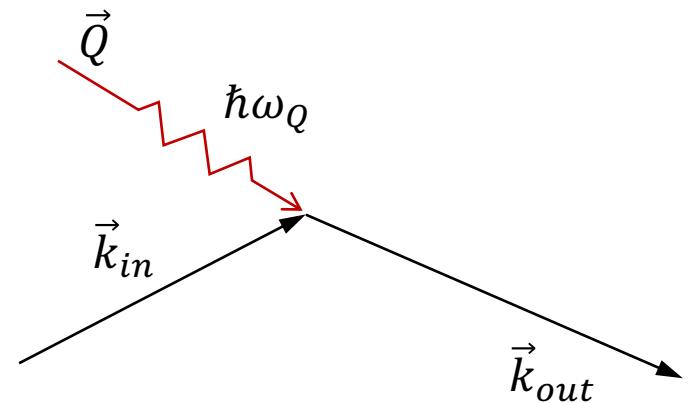


Stokes scattering

$$\frac{\hbar^2 k_{in}^2}{2m_n} = \frac{\hbar^2 k_{out}^2}{2m_n} + \hbar\omega_Q$$

$$\vec{k}_{in} = \vec{k}_{out} + \vec{Q}$$

N-process ("Normal", within the first Brillouin zone)

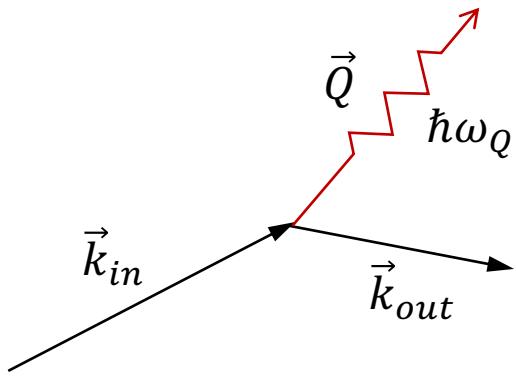


Anti-Stokes scattering

$$\frac{\hbar^2 k_{in}^2}{2m_n} + \hbar\omega_Q = \frac{\hbar^2 k_{out}^2}{2m_n}$$

$$\vec{k}_{in} + \vec{Q} = \vec{k}_{out}$$

# Scattering of neutrons by phonons

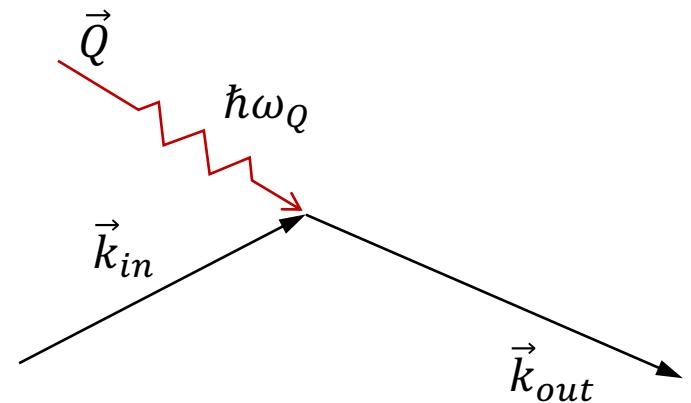


Stokes scattering

$$\frac{\hbar^2 k_{in}^2}{2m_n} = \frac{\hbar^2 k_{out}^2}{2m_n} + \hbar\omega_Q$$

$$\vec{k}_{in} = \vec{k}_{out} + \vec{Q} + \vec{G}$$

U-process (“Umklapp”, outside of the first Brillouin zone)



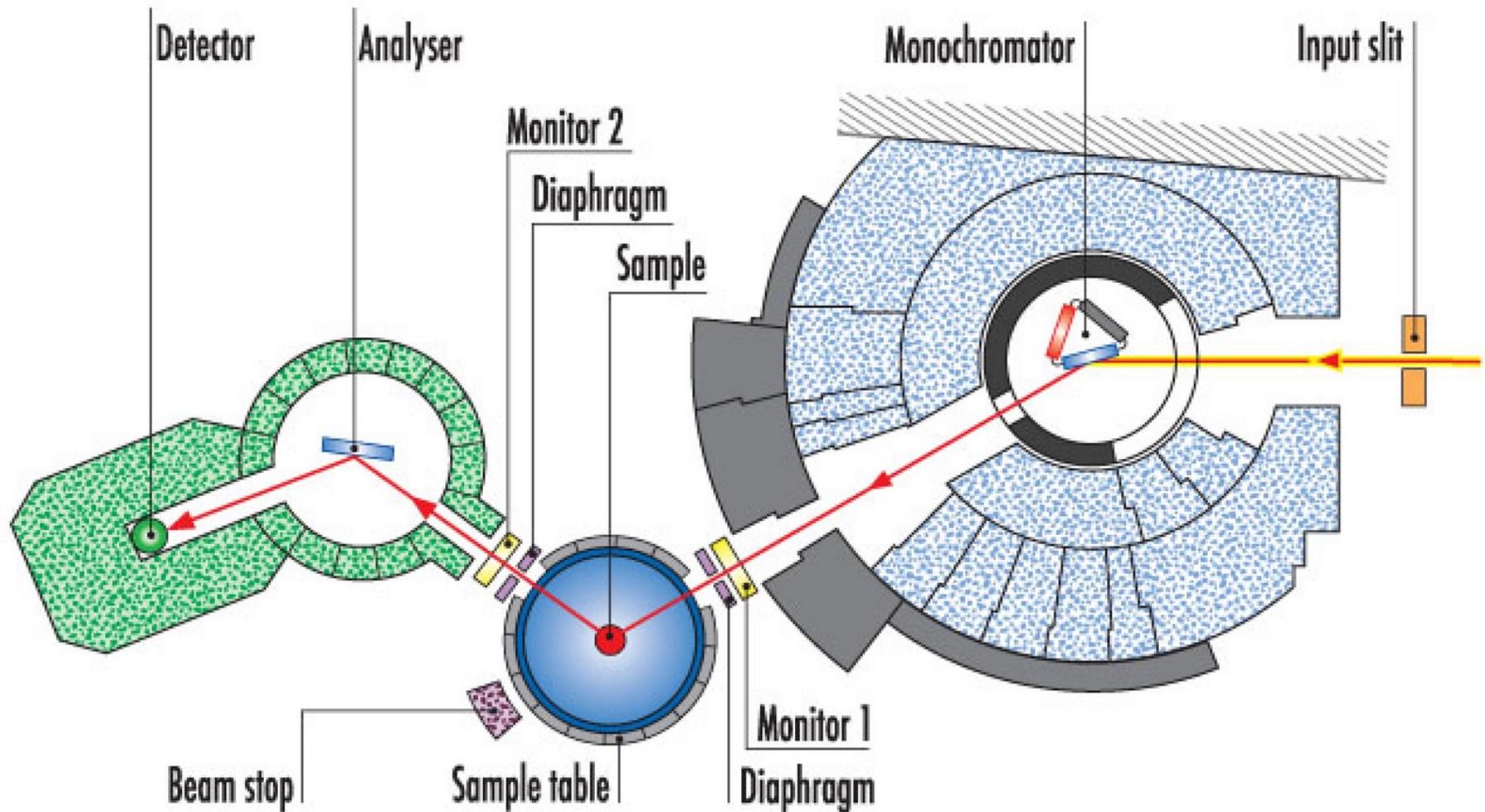
Anti-Stokes scattering

$$\frac{\hbar^2 k_{in}^2}{2m_n} + \hbar\omega_Q = \frac{\hbar^2 k_{out}^2}{2m_n}$$

$$\vec{k}_{in} + \vec{Q} = \vec{k}_{out} + \vec{G}$$

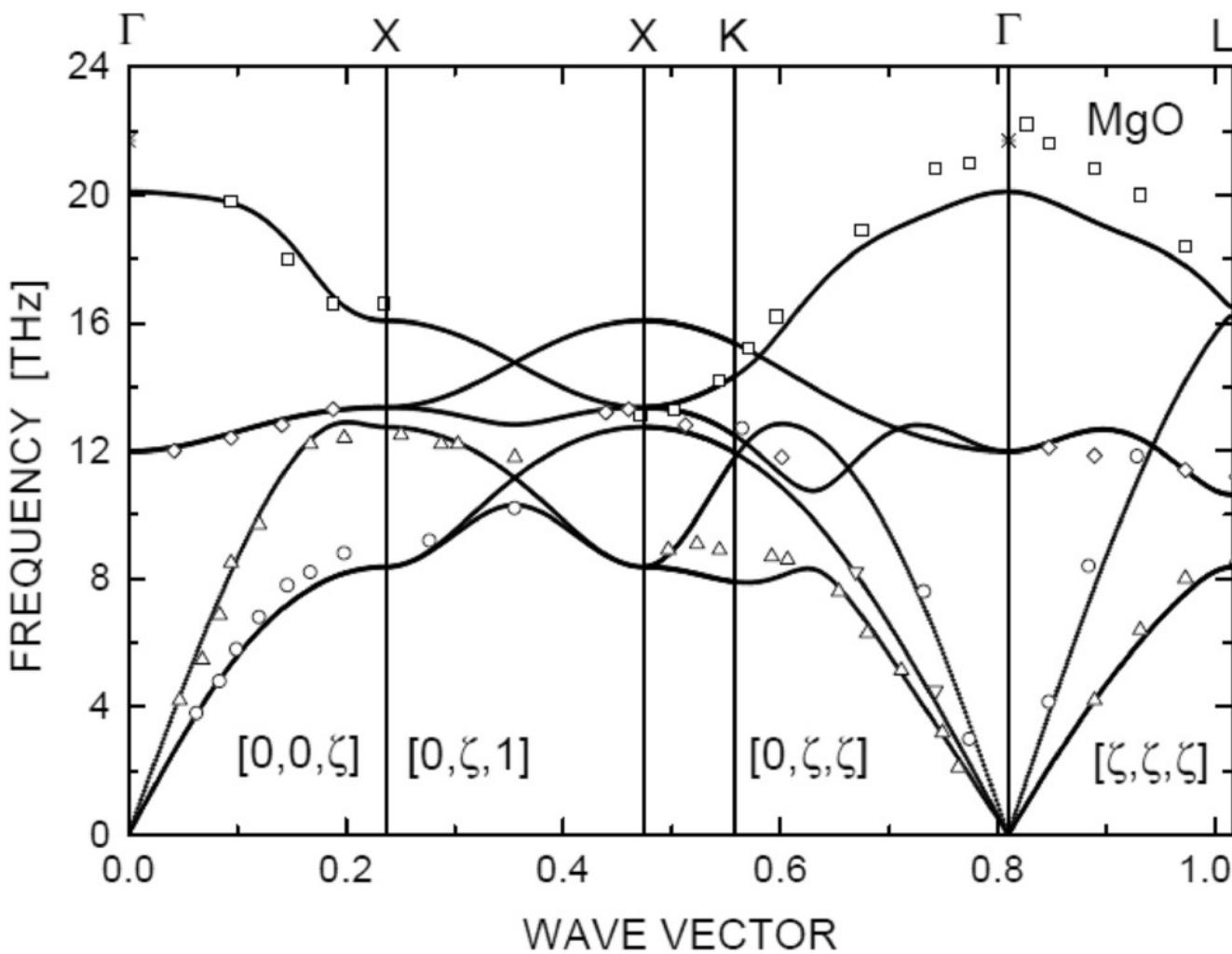
# Measuring dispersion with neutrons

# Experimental setup



*Triple-axis spectrometer: schematic layout of IN8 at a thermal beam (ILL, Grenoble)*

# Experimental setup



K. Parlinski et al., J. Phys. Chem. Solids., **61**, 87–90, (1999)

# Neutrons vs. X-rays

Neutrons:

$$E = \frac{2\pi^2 \hbar^2}{m_n \lambda^2} \quad E[eV] = \frac{0.082}{\lambda^2 [\text{\AA}^2]} \quad \frac{\Delta E}{E} = \frac{\hbar\omega}{E} \sim \frac{0.025 \text{ eV}}{0.082 \text{ eV}} \sim 0.3$$

Easy to measure a 30% change in the energy of a neutron

Photons (x-rays):

$$E = \frac{2\pi\hbar c}{\lambda} \quad E[eV] = \frac{12400}{\lambda[\text{\AA}]} \quad \frac{\Delta E}{E} = \frac{\hbar\omega}{E} \sim \frac{0.025 \text{ eV}}{12400 \text{ eV}} \sim 10^{-6}$$

Change of the X-ray photon energy is negligible. Elastic scattering is a good approximation.

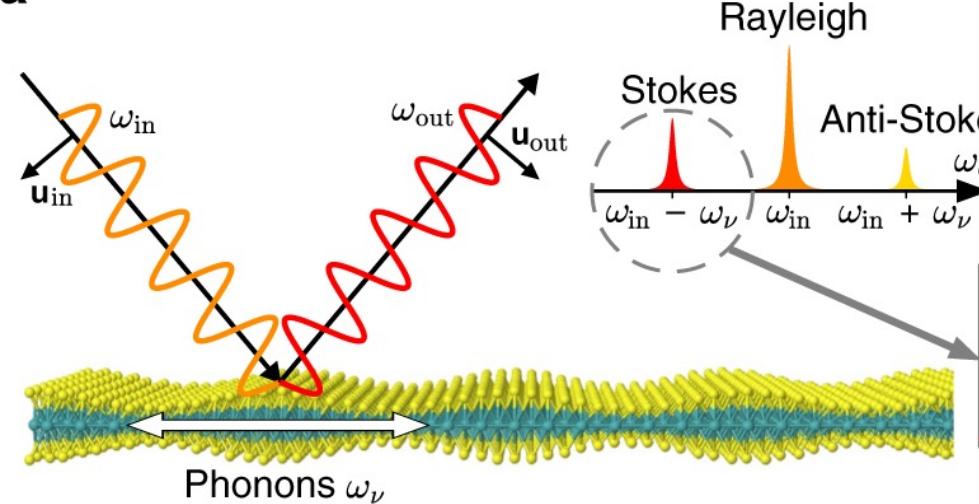
Typical phonon in a crystal:

$$\hbar\omega \sim k_B T \sim 0.025 \text{ eV}$$

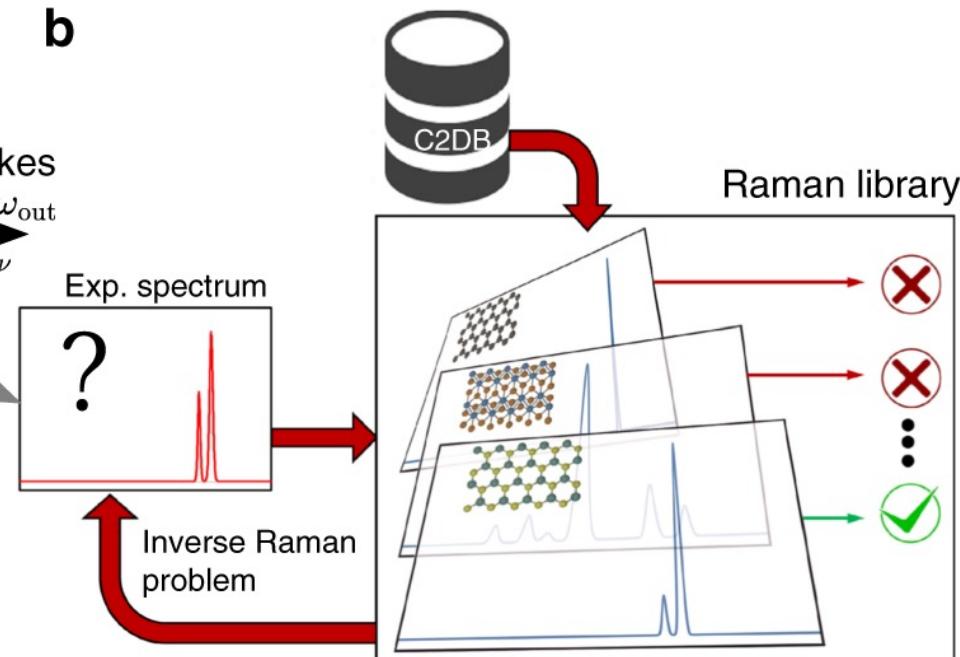
$$q \sim \frac{2\pi}{a} \sim 2 \text{ \AA}^{-1}$$

# Raman scattering

a

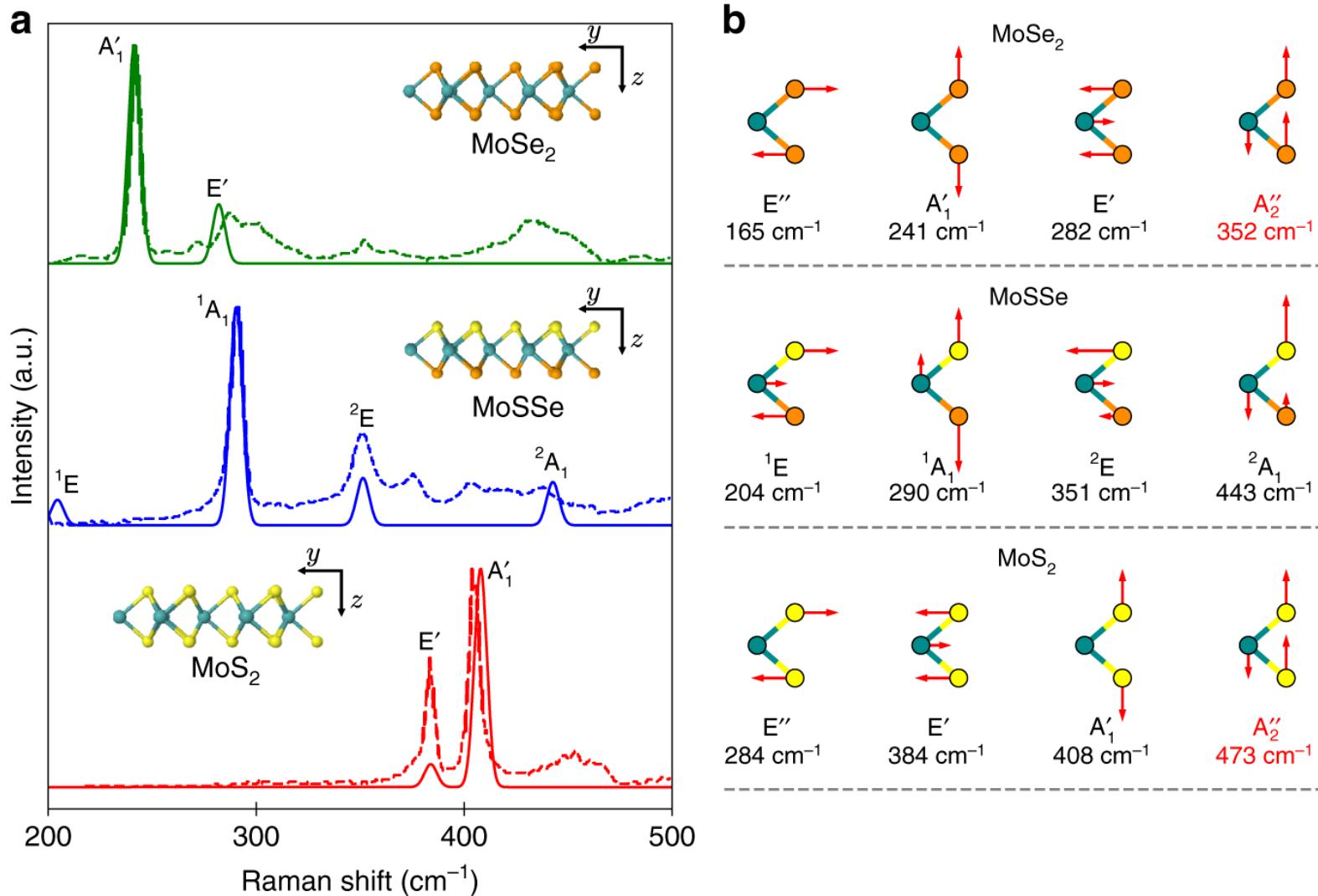


b



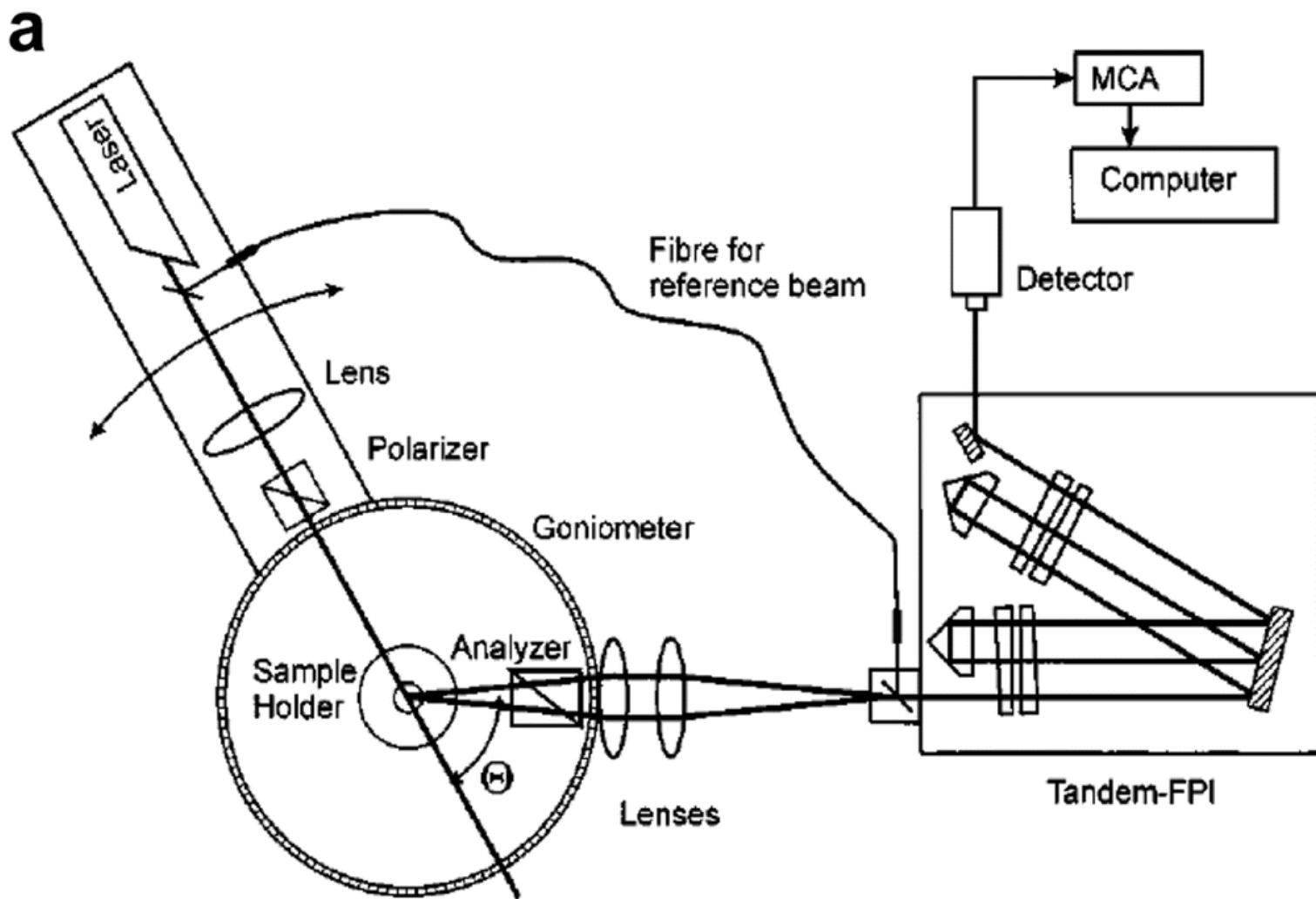
A. Taghizadeh et al., Nat. Commun., 11, 3011, (2020)

# Raman scattering



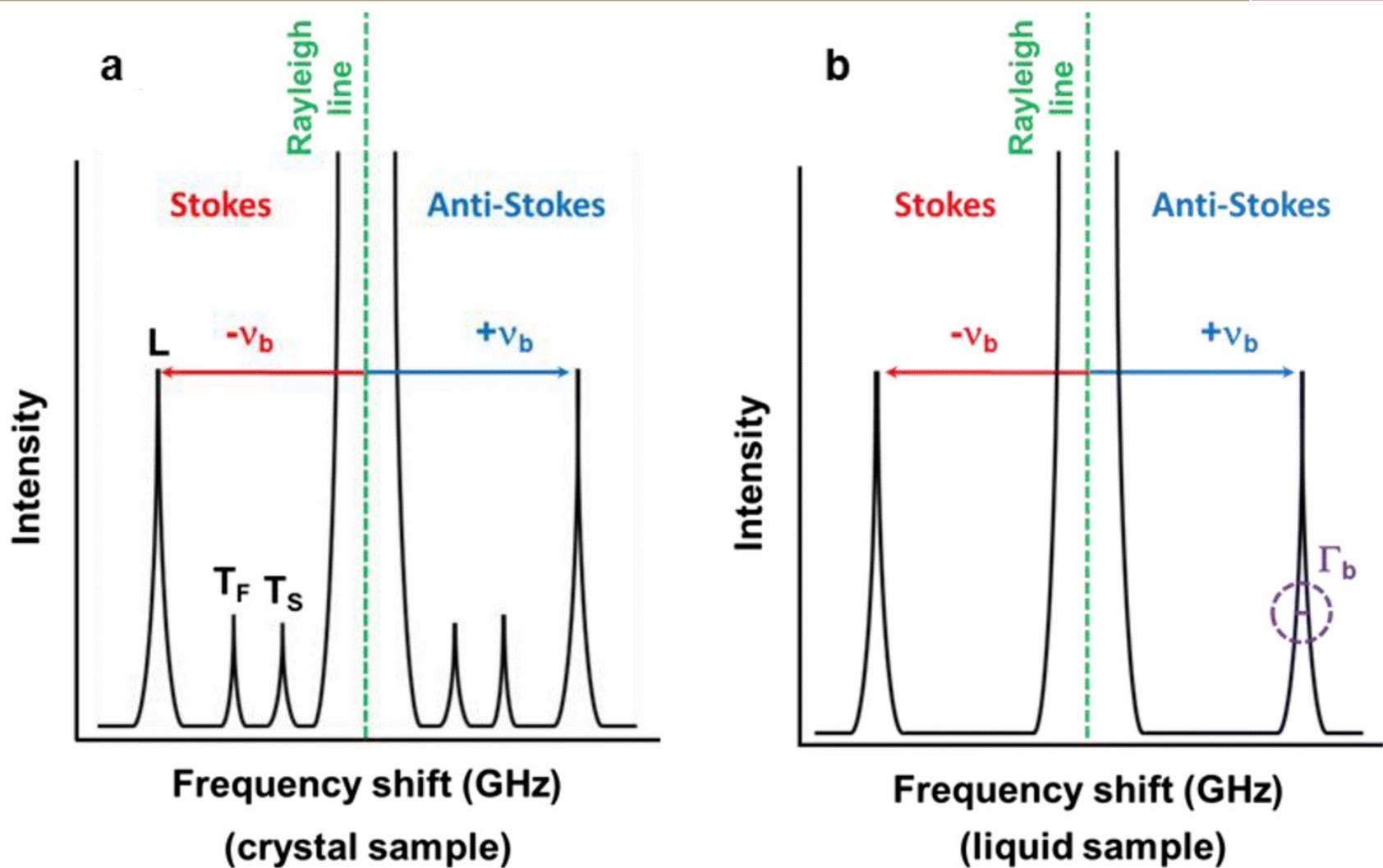
A. Taghizadeh et al., Nat. Commun., 11, 3011, (2020)

# Brillouin scattering



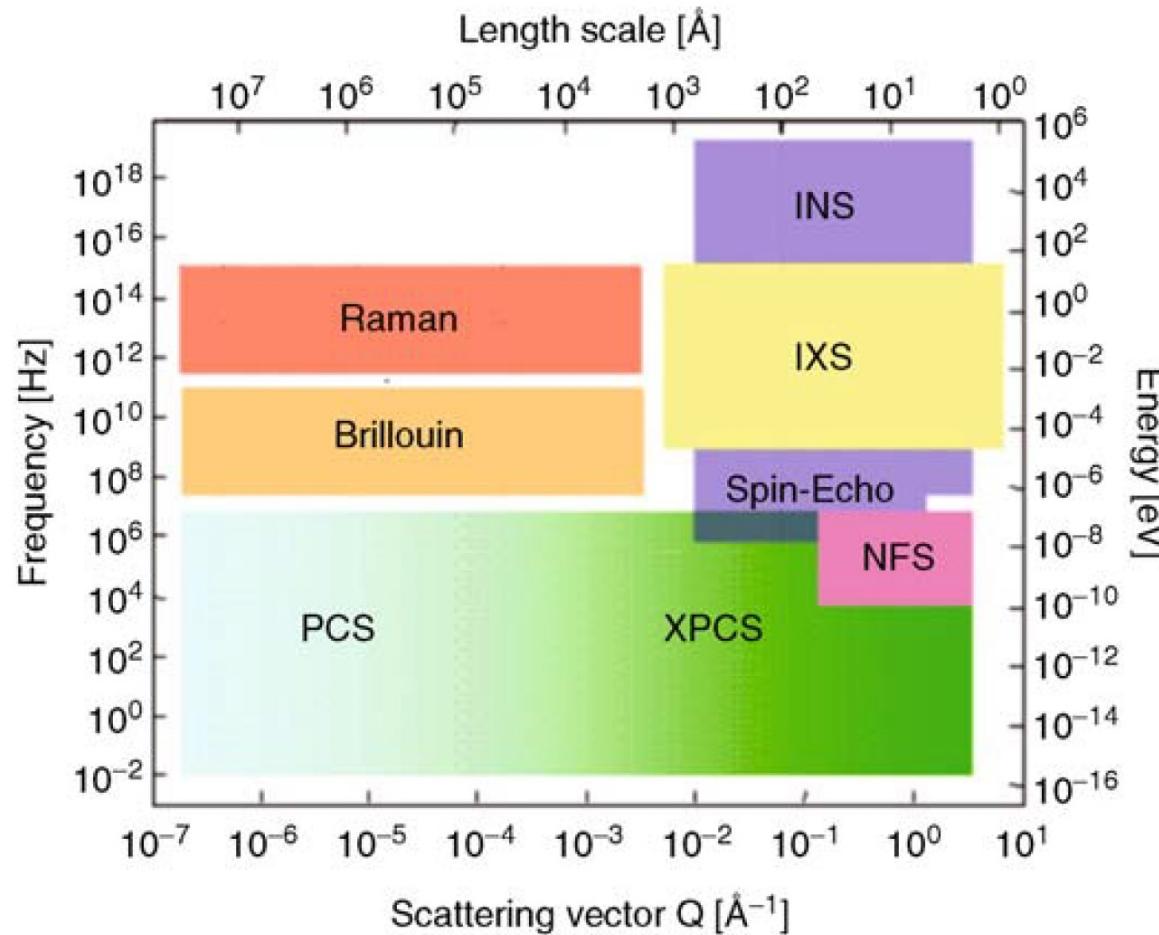
A. Singaraju et al., AAPS PharmSciTech., 20, 109, (2019)

# Brillouin scattering



A. Singaraju et al., AAPS PharmSciTech., 20, 109, (2019)

# Time-resolved techniques



Raman – inelastic scattering of optical light by molecules  
Brillouin – inelastic scattering of optical light by phonons  
PCS – photon correlation spectroscopy  
XPCS – x-ray photon correlation spectroscopy

NFS – nuclear forward scattering  
NSE – neutron spin-echo spectroscopy  
IXS – inelastic x-ray scattering  
INS – inelastic neutron scattering

# What to remember

- Elastic scattering is sensitive to the time-averaged structure

$$G(\vec{r}) = \frac{1}{N} \left\langle \sum_{n,m=1}^N \delta \left( r - ((\vec{r}_n - \vec{r}_m)) \right) \right\rangle$$

$$S(\vec{q}) = \int G(\vec{r}) e^{-i\vec{q}\vec{r}} d\vec{r}$$

- Inelastic scattering is sensitive to the dynamics

$$G(\vec{r}, t) = \frac{1}{N} \left\langle \sum_{n,m=1}^N \delta \left( r - ((\vec{r}_n(t) - \vec{r}_m(0))) \right) \right\rangle$$

$$S_{coh}(\vec{q}, \omega) = \frac{1}{2\pi\hbar} \int G(\vec{r}, t) e^{-i\vec{q}\vec{r}} e^{i\omega t} d\vec{r} dt$$

- Inelastic neutron scattering is the best tool to measure phonon dispersion
- Raman scattering involves optical phonons
- Brillouin scattering involves acoustic phonons